ANSWERS

Accelerating Frame and Rotating Disk

Exercise 1 Two rockets with the same proper acceleration synchronized

Equations of the trajectory of an objet in an accelerating frame, at T=0, X=0 and Y=0:

$$X = \frac{1}{\cosh T - \beta_0 \sinh T \cos \theta} - 1$$
 and $Y = \frac{\beta_0 \sinh T \sin \theta}{\cosh T - \beta_0 \sinh T \cos \theta}$,

 β_0 c is the initial speed and θ the angle with the vertical direction. $T=t/t_H$, $X=x/d_H$ and $Y=y/d_H$.

1) Two rockets in the accelerating frame, at rest at A and B, exchange photons to verify they are well synchronized (see the figure, $X_A = X_B = 0$, $Y_A = 0$ and $Y_B = 2$).

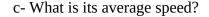
a- How long it takes for a photon to go from *A* to *B*?

The arc AB is a quarter circle $\Rightarrow \theta = \pi/4$. Light speed for X=0: $\beta_0 = 1$.

so
$$X=0 \Rightarrow chT - shT/\sqrt{2} = 1$$
 and $Y=2 \Rightarrow shT = 2\sqrt{2}$
and $T = argsh(2\sqrt{2}) \approx 1.76$

b- What distance did the photon travel?

Circle radius: $R=\sqrt{2}$. Arc $AB=\pi R/2=\pi/\sqrt{2}\simeq 2.22$



 $v/c=\pi/\sqrt{2/argsh} \ 2\sqrt{2} \approx 1.26$

2) Let's now consider the point of view of an inertial observer in R'. At T'=0, T=0 and X'=X=0.

Change of coordinates:

$$T'=(X+1) shT$$
, $X'=(X+1) chT-1$ and $Y'=Y$.

- a- Drawn the two rockets and the ray of light in R'.
- b- Give the coordinates of the events: the photon starts from *A*, and, the photon arrives at *B*.

A: T=0, T'=0, X'=0, Y'=0.

B:
$$T = \operatorname{argsh2}\sqrt{2}$$
, $T' = (0+1) 2\sqrt{2} = 2\sqrt{2}$,

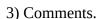
 $X'=(0+1)ch(argsh2\sqrt{2})-1=2$ and Y'=Y=2.

c- How long it takes for a photon to go from *A* to *B*? $\Delta T'=T'_B-T'_A=2\sqrt{2}$

d- What distance did the photon travel?

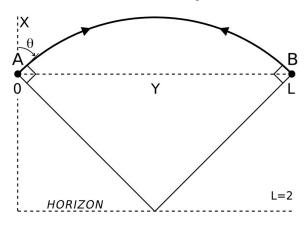
 $AB = \sqrt{(2^2 + 2^2)} = 2\sqrt{2}$

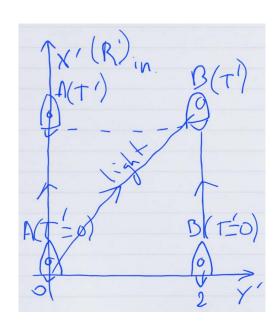
e- What is its average speed? $D'/T'=1=v/c \Rightarrow v=c$



In the non-inertial frame of the rocket, as seen with the metric with $ds^2=0$, the coordinate speed of light is greater than c when X>1 then the average speed v/c is greater than 1.

In an inertial frame of reference we always have the speed of light equal c, and the geodesic is rectilinear.





Exercise 2 Ehrenfest Paradox

 $ds^{2} = \left(1 - \frac{\rho^{2} \omega^{2}}{c^{2}}\right) c^{2} dt^{2} - 2 \rho^{2} \omega dt d\theta - d\rho^{2} - \rho^{2} d\theta^{2} - dz^{2}$ Metric on the rotating disk:

1) Find the expression of the connections $\Gamma^{\lambda}_{\mu\nu}$.

Done on page 412 of the book SRfriends.pdf:

$$g_{tt} = g_{00} = 1 - \frac{\omega^2 \rho^2}{c^2}$$
, $g_{t\theta} = g_{02} = g_{20} = -2\omega \frac{\rho^2}{c}$, $g_{11} = -1$, $g_{22} = -\rho^2$ and $g_{33} = -1$.

The metric is not diagonal. Only the
$$\partial_1 g_{00}$$
, $\partial_1 g_{22}$ and $\partial_1 g_{02}$ are different from zero. Also $g^{\mu\nu}g_{\nu\sigma}=\delta^{\mu}_{\sigma}\Rightarrow g^{33}=1/g_{33}$, $g^{11}=1/g_{11}$, but, $g^{02}g_{20}+g^{00}g_{00}=1$, $g^{02}g_{22}+g^{00}g_{02}=0$, $g^{22}g_{22}+g^{20}g_{02}=1$ and $g^{22}g_{20}+g^{20}g_{00}=0$

then:
$$g^{\mu\nu} = \begin{pmatrix} 1 & 0 & -\frac{\omega}{c} & 0 \\ 0 & -1 & 0 & 0 \\ -\frac{\omega}{c} & 0 & \frac{\omega^2}{c^2} - \frac{1}{\rho^2} & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$
. $\partial_1 g_{00} = -2\frac{\omega^2 \rho}{c^2}, \, \partial_1 g_{22} = -2\rho, \, \partial_1 g_{02} = -4\frac{\omega \rho}{c}.$

Connections symmetric on the covariant indices.

$$\Gamma^{2}_{12} = \frac{1}{2}g^{22}(\partial_{1}g_{22} + \partial_{2}g_{21} - \partial_{2}g_{12}) + \frac{1}{2}g^{20}(\partial_{1}g_{02} + \partial_{2}g_{01} - \partial_{0}g_{12})$$

 $\Gamma^{1}_{00} = \frac{1}{2}g^{11}(\partial_{0}g_{10} + \partial_{0}g_{10} - \partial_{1}g_{00})$ and after calculations only 5 non null connections:

$$\Gamma^{1}_{00} = -\frac{\rho \omega^{2}}{c^{2}} \qquad \Gamma^{1}_{02} = -\frac{\rho \omega}{c} \qquad \Gamma^{1}_{22} = -\rho \qquad \Gamma^{2}_{10} = \frac{\omega}{\rho c} \qquad \Gamma^{2}_{12} = \frac{1}{\rho}$$

2) Find the components of the Riemann tensor $R^{\alpha}_{\beta\mu\nu}$. Comments.

20 independent components:

$$R_{001}^{1} = 0 - \Gamma_{00,1}^{1} + \Gamma_{\sigma 0}^{1} \Gamma_{01}^{\sigma} - 0 = \frac{\omega^{2}}{c^{2}} - \rho \frac{\omega}{c} \frac{\omega}{\rho c} = 0$$

... etc, done page 415 ... all the components are null, spacetime is flat.

Normal, because the metric is obtain from the minkowskian metric with a global change of coordinates: $R^{\alpha}_{\beta\mu\nu} = \frac{\partial x^{\alpha}}{\partial x^{\prime\alpha'}} \frac{\partial x^{\prime\beta'}}{\partial x^{\beta}} \frac{\partial x^{\prime\mu'}}{\partial x^{\mu}} \frac{\partial x^{\prime\nu'}}{\partial x^{\nu}} (R^{\prime\alpha'}_{\beta'\mu'\nu'})_{Mink} = 0$

3) Now, let's find the space curvature.

a- When the disk is motionless, give the metric and the surface of the disk of radius *R*.

$$ds^2 = c^2 dt^2 - d\rho^2 - \rho^2 d\theta^2 - dz^2$$
 (Minkowski), Euclidean space: $S = \pi R^2$.

b- When the disk rotates with an angular speed ω the reference system, has said by *Landau*, is not "synchronous", $g_{0i} \neq 0$, the temporal coordinate is not directly separated from the spatial coordinates. Then, it is shown that the spatial metric is:

$$dl^2 = \gamma_{ij} dx^i dx^j$$
 with $\gamma_{ij} = -g_{ij} + \frac{g_{0i}g_{0j}}{g_{00}}$

(Landau/Lifchitz, The Classical Theory of Field, § Distances and time intervals)

The lengths and the surface $S = \int \sqrt{|\det y_{ii}|} dx^i dx^j$ can be calculate:

i) Express the three-dimensional metric tensor γ_{ii}.

Page 419: ...
$$\gamma_{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \gamma^2 \rho^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

ii) Determine the diameter D and the perimeter P of the disk. What is the value of P/D? Comments.

Page 420: ...
$$\frac{P}{D} = \frac{\int\limits_{\theta=0}^{\theta=2\pi} \sqrt{\gamma_{22}} d\theta}{\int\limits_{\rho=0}^{\rho} \sqrt{\gamma_{11}} d\rho} = \frac{\gamma \rho \int\limits_{\theta=0}^{\theta=2\pi} d\theta}{\int\limits_{\rho=0}^{\rho} \sqrt{\gamma_{11}} d\rho} = \gamma \pi \times \pi, \text{ space is curve (even if spacetime is flat).}$$
After, an other way to prove that the space of the disk is curve,

After, an other way to prove that the space of the disk is curve, also in the book, the calculation of R^1_{212} is done: $R^1_{212} = -3 \beta^2 \gamma^6 \neq 0$, and K < 0.

iii) Find S and compare with the euclidean surface when $\omega \ll c/R$.

$$\det \gamma_{ij} = 1 \times \gamma^{2} \rho^{2} \times 1 \text{ and } \sqrt{|\det \gamma_{ij}|} = \frac{\rho}{\sqrt{1 - \frac{\rho^{2} \omega^{2}}{c^{2}}}} \text{ then } S = \int_{\theta=0}^{2\pi} \int_{\rho=0}^{R} \frac{\rho \, d\rho \, d\theta}{\sqrt{1 - \frac{\rho^{2} \omega^{2}}{c^{2}}}} = 2\pi \frac{c^{2}}{\omega^{2}} \int_{u=0}^{R^{\frac{\omega}{c}}} \frac{u \, du}{\sqrt{1 - u^{2}}}$$

$$\dots S = 2\pi \frac{c^{2}}{\omega^{2}} \left[1 - \sqrt{1 - \left(R \frac{\omega}{c} \right)^{2}} \right] \simeq \pi R^{2} + \frac{1}{8} \left(\frac{R \omega}{c} \right)^{4} > S_{Euclid}$$